Generation and O\textsubscript{2} Adsorption Studies of the Microporous Magnets CsNi[Cr(CN)\textsubscript{6}] (T\textsubscript{C} = 75 K) and Cr\textsubscript{3}[Cr(CN)\textsubscript{6}]\textsubscript{2}·6H\textsubscript{2}O (T\textsubscript{N} = 219 K)

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Abstract: Dehydration of the Prussian blue analogues CsNi[Cr(CN)\textsubscript{6}]-2H\textsubscript{2}O (1) and Cr\textsubscript{3}[Cr(CN)\textsubscript{6}]\textsubscript{2}-10H\textsubscript{2}O (2) affords two new microporous magnets: CsNi[Cr(CN)\textsubscript{6}] (1d) and Cr\textsubscript{3}[Cr(CN)\textsubscript{6}]\textsubscript{2}-6H\textsubscript{2}O (2d). Compounds 1d and 2d maintain the Prussian blue structure, and N\textsubscript{2} adsorption measurements at 77 K show them to be microporous with BET surface areas of 360 and 400 m\textsuperscript{2}/g, respectively. Both solids largely retain the magnetic properties of their parent hydrates, with 1d ordering at 75 K and 2d ordering at 219 K, by far the highest ordering temperature yet observed for a microporous magnet. The compounds further show unexpected changes in their magnetic properties upon adsorption of O\textsubscript{2}. In 2d, adsorption of O\textsubscript{2} results in a reversible decrease in the magnetic moment of the system, as well as a reduction of the coercivity from 110 to 10 G and of the remnant magnetization from 1200 to 400 emu/G/mol, indicating a net antiferromagnetic interaction between O\textsubscript{2} and the framework. In 1d, adsorption of O\textsubscript{2} instead results in a reversible increase in the magnetic moment of the system, indicating a net ferromagnetic interaction between O\textsubscript{2} and the framework. Together, the results suggest that ferromagnetic exchange coupling between O\textsubscript{2} and the [Cr(CN)\textsubscript{6}]\textsuperscript{3-} units provides the predominant magnetic interaction of the adsorbate with the framework.

Introduction

The reduced number of bonds per unit volume associated with a microporous solid is at direct odds with the correlation between the number of exchange pathways and the ordering temperature of a magnet.\textsuperscript{1} In addition, the synthesis of a microporous solid that behaves as a magnet at room temperature remains an open challenge. Such compounds are of interest for their potential utility as low-density magnets or magnetic sensors, and possibly even for use in performing magnetic separations, in which the magnetic flux within the pores of the material would selectively attract their magnetic properties, with 1d ordering at 75 K and 2d ordering at 219 K, by far the highest ordering temperature yet observed for a microporous magnet. The compounds further show unexpected changes in their magnetic properties upon adsorption of O\textsubscript{2}. In 2d, adsorption of O\textsubscript{2} results in a reversible decrease in the magnetic moment of the system, as well as a reduction of the coercivity from 110 to 10 G and of the remnant magnetization from 1200 to 400 emu/G/mol, indicating a net antiferromagnetic interaction between O\textsubscript{2} and the framework. In 1d, adsorption of O\textsubscript{2} instead results in a reversible increase in the magnetic moment of the system, indicating a net ferromagnetic interaction between O\textsubscript{2} and the framework. Together, the results suggest that ferromagnetic exchange coupling between O\textsubscript{2} and the [Cr(CN)\textsubscript{6}]\textsuperscript{3-} units provides the predominant magnetic interaction of the adsorbate with the framework.

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temperatures can be constructed. Prominent among these are CsNi[Cr(CN)6]·2H2O (Tc = 90 K), Cr3[Cr(CN)6]2·10H2O (Tc = 240 K), V[Cr(CN)6]3+ (Tc = 315 K), and KV[Cr(CN)6]2 (Tc = 376 K). Prussian blue analogues have also shown the potential for high surface area, with dehydrated compounds of the type M2[Fe(CN)6]4− and M3[Fe(CN)6]2− exhibiting BET surface areas ranging from 520 to 870 m2/g. To date, however, the only Prussian blue analogues in which microporosity and bulk magnetic ordering have been demonstrated to coexist are M3[Fe(CN)6]2− (Tc = 10 K for M = Mn, 14 K for M = Co, 25 K for M = Ni, and 22 K for M = Cu).10 Cr3[Cr(CN)6]2− (Tc = 38 K), and K2O·Mn4[Cr(CN)6] (Tc = 99 K). Herein, we report desolvation of the Prussian blue analogues CsNi[Cr(CN)6]·2H2O (1) and Cr3[Cr(CN)6]2·10H2O (2) to give CsNi[Cr(CN)6]·1H2O (1d) and Cr3[Cr(CN)6]2·6H2O (2d), microporous magnets with BET surface areas of 360 and 400 m2/g and magnetic ordering temperatures of 75 and 219 K, respectively. Unexpectedly, both materials were found to exhibit changes in their magnetic properties upon adsorption of O2 but not N2, indicating the presence of magnetic exchange coupling between the O2 guest molecules and the frameworks.

Experimental Section

Preparation of Compounds. Reactions involving Cr2+ salts were performed inside a glovebag under a nitrogen atmosphere. The compound K1[Cr(CN)6] was prepared according to a previously reported procedure. Water was distilled and deionized with a Milli-Q filtering system and was subsequently degassed by rapidly stirring under reduced pressure (ca. 10−3 bar) for 4 h. All other reagents were obtained from commercial vendors and, unless otherwise noted, were used without further purification.

CsNi[Cr(CN)6]·2H2O (1). The synthesis of this compound was based upon a modification of a previously reported procedure. A solution of NiCl2 (240 mg, 1.0 mmol) in 30 mL of H2O was added dropwise to a stirred solution of K1[Cr(CN)6] (330 mg, 1.0 mmol) and CsCl (340 mg, 2.0 mmol) in 10 mL of water. The resulting light blue precipitate was allowed to anneal in the mother liquor for 1 h and then was collected by filtration and washed with successive aliquots of water (3 × 30 mL). The solid was dried under reduced pressure to afford 400 mg (92%) of product as a yellow microcrystalline powder. Anal. Calcd for CsH8CrCsNiO4: C, 16.54; H, 0.93; Cr, 11.93; N, 19.29; Ni, 13.47. Found: C, 16.85; H, 1.17; Cr, 11.99; N, 19.57; Ni, 13.24. Elemental analysis further indicated the presence of less than 0.05 equiv of K. The powder X-ray diffraction pattern of this compound agrees with that of a typical Prussian blue type solid exhibiting a face-centered cubic unit cell.

CsNi[Cr(CN)6] (1d). A sample of 1 (100 mg, 0.75 mmol) was heated to 95 °C under reduced pressure (ca. 10−4 bar) for 48 h to give the product as a yellow microcrystalline solid in quantitative yield. Anal. Calcd for CsH8CrNi3O: C, 18.03; Cr, 13.01; N, 21.02; Ni, 14.68. Found: C, 18.24; Cr, 12.87; N, 21.24; Ni, 14.79. Elemental analysis further indicated the presence of less than 0.05 equiv of K. The powder X-ray diffraction pattern of this compound agrees with that of a typical Prussian blue type solid exhibiting a face-centered cubic unit cell.

Cr3[Cr(CN)6]2·10H2O (2). The synthesis of this compound was based upon a modification of a previously reported procedure. A solution of CrCl2 (370 mg, 3.0 mmol) in 10 mL of water was filtered through a plug of Celite to remove insoluble impurities. A solution of K1[Cr(CN)6] (330 mg, 1.0 mmol) in 10 mL of water was then added dropwise to the purified CrCl2 solution under vigorous stirring. The resulting gray precipitate was allowed to anneal in the mother liquor for 12 h and was then collected by filtration and washed with successive aliquots of water (3 × 30 mL). The solid was dried in air to afford 360 mg (96%) of product as a gray microcrystalline powder. Anal. Calcd for Cs2H25Cr3N12O12: C, 19.16; H, 2.68; Cr, 34.6; N, 22.34. Found: C, 19.47; H, 2.55; Cr, 34.1; N, 22.20. Elemental analysis further indicated the presence of less than 0.05 equiv of K. The powder X-ray diffraction pattern of this compound agrees with that of a typical Prussian blue type solid exhibiting a face-centered cubic unit cell.

Cr3[Cr(CN)6]2·6H2O (2d). A sample of 2 (100 mg, 0.13 mmol) was heated to 85 °C under reduced pressure (ca. 10−4 bar) for 48 h to give the product as a gray microcrystalline solid in quantitative yield. Anal. Calcd for Cs2H25Cr3N12O12: C, 21.19; H, 1.78; Cr, 38.22; N, 24.71. Found: C, 21.36; H, 1.63; Cr, 38.12; N, 24.60. Elemental analysis further indicated the presence of less than 0.05 equiv of K. The powder X-ray diffraction pattern of this compound agrees with that of a typical Prussian blue type solid exhibiting a face-centered cubic unit cell.

Gas Adsorption Measurements. Samples of known weight were loaded with 100−300 mg of sample and sealed using a transcale. Samples were degassed at 85 or 95 °C for 24−48 h on a Micromeritics ASAP 2020 analyzer until the outgas rate was less than 1 mTorr/min. The degassed sample and sample tube were weighed precisely and then transferred back to the analyzer (with the transcale preventing exposure of the sample to air after degassing). The outgas rate was again confirmed to be less than 1 mTorr/min. Samples were maintained at constant temperature by immersion in baths of liquid nitrogen (77 K), liquid argon (87 K), ethyl alcohol/liquid nitrogen (157 K), heptane/liquid nitrogen (182 K), acetone/dry ice (195 K), octane/liquid nitrogen (217 K), or acetonitrile/liquid nitrogen (232 K). UHP grade N2, O2, and He (99.999%) gases were used for all measurements.

Results and Discussion

Enthalpies of adsorption were calculated by fitting the data to the van’t Hoff equation:

$$\ln(P) = \frac{\Delta H_{\text{ads}}}{RT} + \frac{\Delta S_{\text{ads}}}{R}$$

(1)

where $P$ is the pressure, $T$ is the temperature, $R$ is the molar gas constant, $\Delta H_{\text{ads}}$ is enthalpy of adsorption, and $\Delta S_{\text{ads}}$ is the entropy of adsorption. This equation can be used to calculate the enthalpy of adsorption as a function of the quantity of gas adsorbed. Pressure as a function of the quantity of gas adsorbed was calculated by fitting each isotherm using the Langmuir–Freundlich equation:

$$Q = \frac{Q_m B \cdot P^{1/\alpha}}{1 + B \cdot P^{1/\alpha}}$$

(2)

where $Q$ is the number of moles of gas adsorbed, $Q_m$ is the number of moles of gas adsorbed at saturation, $P$ is the pressure, and $B$ and $t$ are fitting constants. In order to obtain the most accurate interpolation between measured data points, fits were only applied to regions of the isotherm that were used to calculate the enthalpy of adsorption. Similar results were obtained when virial equations were employed in fitting the isotherms.22

Magnetic Susceptibility Measurements. Quartz sample tubes were loaded with 15–30 mg of sample and degassed under reduced pressure (10−7 bar) at 85 or 95 °C for 48 h. The tubes were then flame-sealed, and the magnetic susceptibility measurements were performed using a Quantum Design MPMS2 SQUID magnetometer. For samples measured in the presence of adsorbed gas, prior to sealing the sample tube, the degassed samples were cooled to 77 K and the sample tubes pressurized to 50 Torr with dioxygen or dinitrogen for 10 min. The amount of gas adsorbed by the sample was calculated using the gas adsorption measurements described above.

Other Physical Measurements. Elemental analyses for C, H, and N were obtained from the Microanalytical Laboratory of the University of California, Berkeley, and analyses for K, Cr, and Ni were obtained from Huffman Laboratories. Thermogravimetric analyses were carried out at a ramp rate of 0.5 °C/min using a TA Instruments Q5000 analyzer. Powder X-ray diffraction data were collected using Cu Kα ($\lambda = 1.5406$ Å) radiation on a Siemens D5000 diffractometer.

Figure 2. Nitrogen adsorption isotherms for 1d (red squares) and 2d (blue circles), as measured at 77 K.

Figure 3. Field-cooled magnetization data for 1d (red) and 2d (blue), as measured in an applied field of 1000 Oe.

50 Torr due to adsorption in the internal micropores, followed by a gradual increase from 50–660 Torr, likely due to adsorption on the surface of the particles. Fits of the data to the BET equation gave surface areas of 360 and 400 m²/g for 1d and 2d, respectively. The lower surface areas of these materials relative to the 520–870 m²/g reported for a variety of other Prussian blue analogues15–18 are presumably due to the partial occupancy of the pores by Cs+ ions in the case of 1d and H2O in the case of 2d, together with the lack of cyanometallate vacancies in the former case.

Remarkably, the magnetic properties of 1 and 2 are largely retained by the dehydrated materials. As shown in Figure 3, magnetization data for 1d indicate a ferromagnet with an ordering temperature of $T_C = 75$ K, compared with $T_C = 90$ K for 1.10 The room temperature magnetic susceptibility for 1d gives $\chi_M T = 3.32$ emu K/mol, similar to the value of 2.88 emu K/mol expected for one Ni²⁺ (S = 1) ion and one Cr³⁺ (S = 3/2) ion per formula unit, assuming no exchange coupling and $g = 2.00$. The moment of 1d continuously rises with decreasing temperature, and a fit of the data to the Curie–Weiss law gives a positive Weiss constant of $\theta = 74$ K (see Figure S1 in the Supporting Information), indicating the expected ferromagnetic coupling between Ni²⁺ and Cr³⁺ ions. No significant magnetic hysteresis is observed for 1d at temperatures down to 2 K, which is in accordance with the very small coercivity ($H_c = 70$ Oe) observed for the parent compound 1.10 As the magnetic coercivity can also be reduced by smaller grain size, even if Ni²⁺ has been known to possess significant single-ion anisotropy, compound 1d is likely

to lose anisotropy due to the very small crystallite sizes as well as the high symmetry of the Prussian blue structure type.\(^{(23)}\)

The magnetization data for \(2d\), also shown in Figure 3, are indicative of a ferrimagnet with an ordering temperature of \(T_N = 219\) K, nearly identical to the value of \(T_N = 220\) K observed for \(2\).\(^{(24)}\) This represents the highest ordering temperature yet reported for a microporous magnet. The room temperature magnetic susceptibility for \(2d\) gives \(\chi_M T = 5.64\) emu K/mol, somewhat lower than the value of 12.75 emu K/mol expected for two high-spin \(\text{Cr}^{2+}\) (\(S = 2\)) ions and three \(\text{Cr}^{3+}\) (\(S = \frac{5}{2}\)) ions per formula unit, assuming no exchange coupling and \(g = 2.00\), due to the antiferromagnetic coupling at room temperature.\(^{(12)}\) As the sample is cooled, the moment decreases and reaches a minimum at 287 K. This, together with the negative Weiss constant of \(\theta = -570\) K, obtained using Neel’s hyperbolic equation\(^{(25)}\) (see Figure S2 in the Supporting Information), indicates antiferromagnetic coupling between adjacent metal centers. Consistent with ferrimagnetism, magnetic hysteresis is observed at 5 K, with a loop characterized by a coercive field of 110 Oe and a remnant magnetization of 1200 emu-G/mol, comparable to the values of 20 Oe and 1333 emu-G/mol observed for \(2\).\(^{(12)}\)

To investigate the effect of the magnetic field of a porous magnet on a paramagnetic adsorbent, \(O_2\) adsorption isotherms were measured for \(2d\) at temperatures ranging from 77 to 232 K (see Figure S3 in the Supporting Information). For a gas adsorbing on a diamagnetic material, the enthalpy and entropy of adsorption are typically constant as a function of temperature.\(^{(21)}\) According to the van’t Hoff equation (see eq 1), a plot of \(\ln(P)\) versus \(T^{-1}\) should therefore be linear. In contrast, for a paramagnetic gas adsorbing on a ferrimagnetic material, the enthalpy of adsorption is augmented by the enthalpy associated with the alignment of the magnetic moment of the gas molecules with the magnetic field within the pores of the material. Since the strength of this interaction increases with decreasing temperature, the slope of \(\ln(P)\) versus \(T^{-1}\) should increase as the temperature increases. The van’t Hoff plots constructed from the \(O_2\) adsorption isotherms for \(2d\) are depicted in Figure 4. Importantly, the plots show no significant deviation from linearity, with fits to the isosteres having \(R^2\) values between 0.994 and 0.999. The coverage dependent enthalpy of adsorption calculated from these isosteres is shown in Figure 5 and ranges from 20.9 to 19.1 kJ/mol. The standard deviation in the enthalpy is less than 0.5 kJ/mol over the measured range, indicating that any magnetic interaction between \(O_2\) and the framework must be very weak. Thus, the possibility of utilizing the microporous magnet \(2d\) for the efficient separation of \(O_2\) from air looks to be unrealizable. For \(1d\), which has a much lower ordering temperature than \(2d\), we do not expect significant interaction of \(O_2\) with the internal magnetic field of the solid at such temperatures. Here, \(O_2\) adsorption isotherms were measured at 77 and 87 K, revealing an enthalpy of \(O_2\) adsorption of 12.0 kJ/mol, which is nearly half of that of \(2d\) (Figure S6 in the Supporting Information). This effect is likely associated with the structural differences between the two compounds (see Figure 1), which give rise to significantly different pore environments.

In an attempt to better characterize the weak interactions between adsorbed \(O_2\) molecules and the cyanobridged framework, the magnetic properties of \(1d\) and \(2d\) were studied in the presence of an atmosphere of \(O_2\). In \(2d\), the presence of 2.9 molecules of \(O_2\) per formula unit, nearly all of which should be adsorbed at 77 K, causes a decrease in the magnetic moment of the system, as well as a reduction in the coercivity from 110 to 10 Oe and the remnant magnetization from 1200 to 440 emu-G/mol (see Figure 6). Both the decrease in the magnetic moment and coercivity are fully reversible upon desorption of \(O_2\). Furthermore, no changes in the magnetic properties of \(2d\) are observed upon adsorption of the diamagnetic gas \(N_2\), suggesting that the \(O_2\) molecules are actually coupling with the framework, rather than simply inducing a change in the structure of \(2d\). To our knowledge, this is the first report of magnetic exchange between a microporous magnet and a paramagnetic adsorbate. At higher magnetic field, the magnetization for \(2d-O_2\) does not saturate, but linearly increases up to \(M = 2.12\) \(N\beta\) at 60 G.\(^{(26)}\) This linear increase has been observed previously for adsorbed \(O_2\) molecules on a diamagnetic porous solid,\(^{(27)}\) where magnetic interaction between the aligned \(O_2\) molecules resulted in linear magnetization up to 5 T at temperatures below 60 K. Also, the magnetic study on the frozen bulk \(O_2\) phase has revealed a


\(^{(24)}\) In our hands, compound \(2\) exhibited ordering temperature of \(T_N = 220\) K, as estimated by a fit of the susceptibility data to Neel’s hyperbolic equation. This is somewhat lower than the value of \(T_N = 240\) K reported in ref 12.


\(^{(26)}\) In our measurement, the saturation was far from being complete, as consistently observed for the parent compound \(\text{Cr}_2[\text{Cr}_2(\text{CN})_6]_2\cdot10\text{H}_2\text{O}\) in ref 12, where the maximum magnetization observed was 1.4 \(N\beta\) at 10 K and 70 kOe.

linear increase in magnetization upon the increasing magnetic field up to 50 T. Therefore, among the 2.9 molecules O$_2$ per formula unit in 2d-O$_2$, some portion is apparently not participating in the magnetic coupling with the framework due to the small pore size but has magnetic interaction between themselves. The complexity of the integrated magnetic interactions including those between O$_2$ molecules, O$_2$ and the framework, and Cr$^{2+}$ and Cr$^{3+}$ in the framework hampered the quantitative analysis of the spin magnetization. Indeed, consistent with our observations for 2d, the magnetization of the parent compound 2 also does not saturate under the conditions probed. A somewhat larger effect is observed for O$_2$ adsorbed within 1d. Here, adsorption of 1.8 molecules of O$_2$ per formula unit causes an increase in the moment of the system, as shown in Figure 7. As with the 2d-O$_2$ system, this change in the magnetic properties is fully reversible upon desorption of O$_2$, and no effect is observed upon adsorption of N$_2$, indicating that O$_2$ is coupling ferromagnetically with the framework. Although the specifics of the exchange mechanism are unclear, we note that the switch from a net antiferromagnetic effect in 2d-O$_2$ to ferromagnetic coupling in 1d-O$_2$ is at least consistent with the predominate interaction being ferromagnetic exchange between O$_2$ and the [Cr(CN)$_6$]$^{3-}$ units of the frameworks. At 2 K, no significant difference in hysteresis was observed between 1d and 1d-O$_2$, other than the change in magnetization values (see Figure 7). Similar to the situation with 2d-O$_2$, the field-dependent magnetization in 1d-O$_2$ is highly increasing without saturation due to the combination of the ferromagnetic interaction between 1d and O$_2$, and the magnetization increase of O$_2$, thus preventing a quantitative spin analysis.

The foregoing results show that microporous magnets with high ordering temperatures can be generated by dehydrating Prussian blue analogues already known to exhibit strong magnetic exchange. Although the magnetic flux within the pores of 2d is insufficient to cause a significant deviation from van’t Hoff behavior, the adsorption of O$_2$ does result in a net antiferromagnetic interaction with the ferrimagnetic framework. Interestingly, this exchange interaction switches to ferromagnetic for O$_2$ adsorbed within the ferromagnet 1d, suggesting that ferromagnetic coupling between O$_2$ and the [Cr(CN)$_6$]$^{3-}$ units provides the predominate exchange pathway in each material. Future work will include powder X-ray diffraction measurements to identify the O$_2$ binding sites within these structures, which may lend additional insight into the origin of the magnetic coupling. In addition, attempts will be made to further increase the ordering temperatures and pore flux density within microporous magnets, through, for example, the incorporation of species such as [Mo(CN)$_6$]$^{3-}$ into Prussian blue analogues.

Acknowledgment. This research was funded by DoE Grant No. DE-FG03-01ER15257. We thank the National Science Foundation for providing S.S.K. with a predoctoral fellowship, and Prof. S. Kitagawa for helpful discussions.

Supporting Information Available: Additional plots of adsorption isotherms and magnetic data. This material is available free of charge via the Internet at http://pubs.acs.org.

JA803926Y